

Total Energy and its Calculation in DFT

Yang Wang

Pittsburgh Supercomputing Center

Carnegie Mellon University

As a many-electron problem

With Born-Oppenheimer approximation, the wave function for the N electrons in a solid satisfies the following Schrödinger equation (in the atomic units: $\hbar = 1$, $m_e = 1/2$, $e^2 = 2$):

$$\left[-\sum_{\alpha=1}^N \nabla_{\alpha}^2 + \sum_{\substack{\alpha,\beta \\ \alpha \neq \beta}} \frac{1}{|\mathbf{r}_{\alpha} - \mathbf{r}_{\beta}|} + \sum_{\alpha=1}^N V_{\text{ext}}(\mathbf{r}_{\alpha}) \right] \Psi(\mathbf{r}_1, s_1; \mathbf{r}_2, s_2; \dots; \mathbf{r}_N, s_N) = E \Psi(\mathbf{r}_1, s_1; \mathbf{r}_2, s_2; \dots; \mathbf{r}_N, s_N)$$

$$\text{where } V_{\text{ext}}(\mathbf{r}_{\alpha}) = -\sum_n \frac{2Z_n}{|\mathbf{r}_{\alpha} - \mathbf{R}_n|}$$

The electron density with spin s is given by

$$\rho_s(\mathbf{r}) = N \sum_{s_2 \dots s_N} \int_{\infty} |\Psi(\mathbf{r}, s; \mathbf{r}_2, s_2; \dots; \mathbf{r}_N, s_N)|^2 d^3\mathbf{r}_2 d^3\mathbf{r}_3 \dots d^3\mathbf{r}_N.$$

$$\rho(\mathbf{r}) = \rho_{\uparrow}(\mathbf{r}) + \rho_{\downarrow}(\mathbf{r})$$

and the electronic total energy is

$$\text{Total energy of the solid: } E_{\text{T}} = E + \sum_n \sum_{m \neq n} \frac{Z_n Z_m}{|\mathbf{R}_n - \mathbf{R}_m|}$$

$$E = \sum_{s_1, s_2, \dots, s_N} \int_{\infty} \Psi^*(\mathbf{r}_1, s_1; \mathbf{r}_2, s_2; \dots; \mathbf{r}_N, s_N) \left[-\sum_{\alpha=1}^N \nabla_{\alpha}^2 + \sum_{\substack{\alpha,\beta \\ \alpha \neq \beta}} \frac{1}{|\mathbf{r}_{\alpha} - \mathbf{r}_{\beta}|} + \sum_{\alpha=1}^N V_{\text{ext}}(\mathbf{r}_{\alpha}) \right] \Psi(\mathbf{r}_1, s_1; \mathbf{r}_2, s_2; \dots; \mathbf{r}_N, s_N) d^3\mathbf{r}_1 d^3\mathbf{r}_2 \dots d^3\mathbf{r}_N.$$

As a one-electron problem in DFT

At the ground state, the electronic total energy E is a functional of the electron density ρ

$$E = T_0[\rho] + U[\rho] + E_{\text{XC}}[\rho]$$

$$\rho(\mathbf{r}) = \rho_{\uparrow}(\mathbf{r}) + \rho_{\downarrow}(\mathbf{r})$$

The electron density ρ can be calculated by $\rho_s(\mathbf{r}) = \sum_{\alpha} f_T(\varepsilon_{\alpha,s} - \varepsilon_F) \psi_{\alpha,s}^*(\mathbf{r}) \psi_{\alpha,s}(\mathbf{r})$

where the Kohn-Sham orbital $\psi_{\alpha,s}(\mathbf{r})$ is the self-consistent solution of the following one-electron Schrödinger equation, known as the Kohn-Sham equation,

$$\left[-\nabla^2 + V_{\text{H}}(\mathbf{r}) + V_{\text{ext}}(\mathbf{r}) + V_{\text{XC}}^s[\rho(\mathbf{r})] \right] \psi_{\alpha,s}(\mathbf{r}) = \varepsilon_{\alpha,s} \psi_{\alpha,s}(\mathbf{r})$$

$$V_{\text{ext}}(\mathbf{r}) = -\sum_n \frac{2Z_n}{|\mathbf{r} - \mathbf{R}_n|}$$

and α is the orbital index. For periodic system, $\alpha = \{n, \mathbf{k}\}$, with n the band index, and \mathbf{k} the Bloch wave \mathbf{k} -vector. The Hartree potential is

$$V_{\text{H}}(\mathbf{r}) = \int_{\infty} \frac{2\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d^3\mathbf{r}'$$

$$f_T(\varepsilon_{\alpha,s} - \varepsilon_F) = \frac{1}{1 + e^{(\varepsilon_{\alpha,s} - \varepsilon_F)/k_B T}}$$

The many-body problem becomes a one-electron problem with effective potential given by

$$V_{\text{eff}}^s(\mathbf{r}) = V_{\text{H}}(\mathbf{r}) + V_{\text{ext}}(\mathbf{r}) + V_{\text{XC}}^s[\rho(\mathbf{r})]$$

Decomposition of the electronic total energy

$$E = T_0[\rho] + U[\rho] + E_{\text{XC}}[\rho]$$

The kinetic energy term

$$\begin{aligned} T_0[\rho] &= \sum_{s=\uparrow,\downarrow} \sum_{\alpha} f_T(\varepsilon_{\alpha,s} - \varepsilon_F) \int_{\infty} \psi_{\alpha,s}^*(\mathbf{r}) \left[-\nabla^2 \right] \psi_{\alpha,s}(\mathbf{r}) d^3\mathbf{r} \\ &= \sum_{s=\uparrow,\downarrow} \sum_{\alpha} f_T(\varepsilon_{\alpha,s} - \varepsilon_F) \int_{\infty} \psi_{\alpha,s}^*(\mathbf{r}) \left[\varepsilon_{\alpha,s} - V_{\text{eff}}^s(\mathbf{r}) \right] \psi_{\alpha,s}(\mathbf{r}) d^3\mathbf{r} \\ &= \sum_{s=\uparrow,\downarrow} \sum_{\alpha} \left(\int_{\infty} \varepsilon f_T(\varepsilon - \varepsilon_F) \delta(\varepsilon - \varepsilon_{\alpha,s}) d\varepsilon \right) \int_{\infty} |\psi_{\alpha,s}(\mathbf{r})|^2 d^3\mathbf{r} - \sum_{s=\uparrow,\downarrow} \int_{\infty} V_{\text{eff}}^s(\mathbf{r}) \sum_{\alpha} f_T(\varepsilon_{\alpha,s} - \varepsilon_F) |\psi_{\alpha,s}(\mathbf{r})|^2 d^3\mathbf{r} \\ &= \sum_{s=\uparrow,\downarrow} \left[\int_{\infty} \varepsilon f_T(\varepsilon - \varepsilon_F) \rho_s(\varepsilon) d\varepsilon - \int_{\infty} V_{\text{eff}}^s(\mathbf{r}) \rho_s(\mathbf{r}) d^3\mathbf{r} \right] \end{aligned}$$

where the spin-resolved density of states $\rho_s(\varepsilon)$ is given by

$$\rho_s(\varepsilon) = \sum_{\alpha} \delta(\varepsilon - \varepsilon_{\alpha,s}) \int_{\infty} |\psi_{\alpha,s}(\mathbf{r})|^2 d^3\mathbf{r}$$

The Coulomb energy term

$$\begin{aligned} U[\rho] &= \int_{\infty} d^3\mathbf{r} \int_{\infty} \frac{\rho(\mathbf{r})\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d^3\mathbf{r}' \\ &\quad + \int_{\infty} \rho(\mathbf{r}) V_{\text{ext}}(\mathbf{r}) d^3\mathbf{r} \end{aligned}$$

$$\rho_s(\mathbf{r}) = \sum_{\alpha} f_T(\varepsilon_{\alpha,s} - \varepsilon_F) |\psi_{\alpha,s}(\mathbf{r})|^2$$

Exchange-correlation energy

$$E_{\text{XC}}[\rho] \approx \int_{\infty} \varepsilon_{\text{XC}}[\rho(\mathbf{r})] \rho(\mathbf{r}) d^3\mathbf{r}$$

Electrostatic Potential and Band Energy

For the effective potential: $V_{\text{eff}}^s(\mathbf{r}) = V_{\text{H}}(\mathbf{r}) + V_{\text{ext}}(\mathbf{r}) + V_{\text{XC}}^s[\rho(\mathbf{r})]$

the sum of the first two terms is the electrostatic potential

$$V_{\text{eff}}^s(\mathbf{r}) = V_{\text{Coul}}(\mathbf{r}) + V_{\text{XC}}^s[\rho(\mathbf{r})]$$

$$V_{\text{Coul}}(\mathbf{r}) = V_{\text{H}}(\mathbf{r}) + V_{\text{ext}}(\mathbf{r}) = \int_{\infty} \frac{2\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d^3\mathbf{r}' - \sum_n \frac{2Z_n}{|\mathbf{r} - \mathbf{R}_n|}$$

whose calculation will be discussed in a separate lecture.

At $T = 0$, the first electronic total energy term becomes

$$\int_{\infty} \varepsilon f_T(\varepsilon - \varepsilon_F) \rho(\varepsilon) d\varepsilon = \int_{-\infty}^{\varepsilon_F} \varepsilon \rho(\varepsilon) d\varepsilon = \sum_c \varepsilon_c + \int_{\varepsilon_b}^{\varepsilon_F} \varepsilon \rho_v(\varepsilon) d\varepsilon$$

ε_c is the energy eigenvalue of a core state, whose density of states is just $\delta(\varepsilon - \varepsilon_c)$.

ε_b is the bottom of valence band energy.

$\rho_v(\varepsilon)$ is the density of the valence states.

$$\int_{\varepsilon_b}^{\varepsilon_F} \varepsilon \rho_v(\varepsilon) d\varepsilon \text{ is recognized as band energy}$$

Total Electrostatic Energy of the Solid

The electrostatic energy of the electron system is:

$$U[\rho] = \int_{\infty} d^3\mathbf{r} \int_{\infty} \frac{\rho(\mathbf{r})\rho(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|} d^3\mathbf{r}' + \int_{\infty} \rho(\mathbf{r})V_{\text{ext}}(\mathbf{r})d^3\mathbf{r}$$

$$V_{\text{ext}}(\mathbf{r}) = -\sum_n \frac{2Z_n}{|\mathbf{r}-\mathbf{R}_n|}$$

The electrostatic energy of the solid is given by:

$$V_{\text{Coul}}(\mathbf{r}) = \int_{\infty} \frac{2\rho(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|} d^3\mathbf{r}' - \sum_n \frac{2Z_n}{|\mathbf{r}-\mathbf{R}_n|}$$

$$\begin{aligned} U[\rho] + \sum_n \sum_{m \neq n} \frac{Z_n Z_m}{|\mathbf{R}_n - \mathbf{R}_m|} &= \int_{\infty} d^3\mathbf{r} \int_{\infty} \frac{\rho(\mathbf{r})\rho(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|} d^3\mathbf{r}' - \sum_n \int_{\infty} \rho(\mathbf{r}) \frac{2Z_n}{|\mathbf{r}-\mathbf{R}_n|} d^3\mathbf{r} + \sum_n \sum_{m \neq n} \frac{Z_n Z_m}{|\mathbf{R}_n - \mathbf{R}_m|} \\ &= \int_{\infty} \rho(\mathbf{r}) \left[\int_{\infty} \frac{\rho(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|} d^3\mathbf{r}' - \sum_n \frac{Z_n}{|\mathbf{r}-\mathbf{R}_n|} \right] d^3\mathbf{r} + \sum_n Z_n \left[-\int_{\infty} \frac{\rho(\mathbf{r})}{|\mathbf{r}-\mathbf{R}_n|} d^3\mathbf{r} + \sum_{m \neq n} \frac{Z_m}{|\mathbf{R}_n - \mathbf{R}_m|} \right] \\ &= \frac{1}{2} \int_{\infty} \rho(\mathbf{r}) V_{\text{Coul}}(\mathbf{r}) d^3\mathbf{r} - \frac{1}{2} \sum_n Z_n V_{\text{Coul}}(\mathbf{R}_n) \end{aligned}$$

Total Energy of the Solid

Recognizing the results:

$$V_{\text{eff}}^s(\mathbf{r}) = V_{\text{Coul}}(\mathbf{r}) + V_{\text{XC}}^s[\rho(\mathbf{r})]$$

$$V_{\text{Coul}}(\mathbf{r}) = \int_{\infty} \frac{2\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d^3\mathbf{r}' - \sum_n \frac{2Z_n}{|\mathbf{r} - \mathbf{R}_n|}$$

$$T_0[\rho] = \int_{-\infty}^{\varepsilon_F} \varepsilon \rho(\varepsilon) d\varepsilon - \int_{\infty} V_{\text{Coul}}(\mathbf{r}) \rho(\mathbf{r}) d^3\mathbf{r} - \sum_{s=\uparrow, \downarrow} \int_{\infty} V_{\text{XC}}^s(\mathbf{r}) \rho_s(\mathbf{r}) d^3\mathbf{r}$$

$$U[\rho] + \sum_n \sum_{m \neq n} \frac{Z_n Z_m}{|\mathbf{R}_n - \mathbf{R}_m|} = \frac{1}{2} \left[\int_{\infty} \rho(\mathbf{r}) V_{\text{Coul}}(\mathbf{r}) d^3\mathbf{r} - \sum_n Z_n V_{\text{Coul}}(\mathbf{R}_n) \right]$$

The total energy of the solid is therefore given by

$$E_T = T_0[\rho] + U[\rho] + E_{\text{XC}}[\rho] + \sum_n \sum_{m \neq n} \frac{Z_n Z_m}{|\mathbf{R}_n - \mathbf{R}_m|} \quad \boxed{\int_{-\infty}^{\varepsilon_F} \varepsilon \rho(\varepsilon) d\varepsilon = \sum_c \varepsilon_c + \int_{\varepsilon_b}^{\varepsilon_F} \varepsilon \rho_v(\varepsilon) d\varepsilon}$$

$$= \int_{-\infty}^{\varepsilon_F} \varepsilon \rho(\varepsilon) d\varepsilon - \frac{1}{2} \left[\int_{\infty} \rho(\mathbf{r}) V_{\text{Coul}}(\mathbf{r}) d^3\mathbf{r} + \sum_n Z_n V_{\text{Coul}}(\mathbf{R}_n) \right] - \sum_{s=\uparrow, \downarrow} \int_{\infty} V_{\text{XC}}^s(\mathbf{r}) \rho_s(\mathbf{r}) d^3\mathbf{r} + E_{\text{XC}}[\rho]$$

Electron density and Green function

Consider the electron density which is given by

$$\begin{aligned}\rho_s(\mathbf{r}) &= \sum_{\alpha} f_T(\varepsilon_{\alpha,s} - \varepsilon_F) |\psi_{\alpha,s}(\mathbf{r})|^2 = \sum_{\alpha} \left[\int_{-\infty}^{\infty} f_T(\varepsilon - \varepsilon_F) \delta(\varepsilon - \varepsilon_{\alpha,s}) d\varepsilon \right] |\psi_{\alpha,s}(\mathbf{r})|^2 \\ &= \int_{-\infty}^{\infty} f_T(\varepsilon - \varepsilon_F) \left[\sum_{\alpha} \delta(\varepsilon - \varepsilon_{\alpha,s}) |\psi_{\alpha,s}(\mathbf{r})|^2 \right] d\varepsilon\end{aligned}$$

Using relation: $\text{Im} \lim_{\eta \rightarrow 0} \frac{1}{z + i\eta} = -\pi\delta(z)$, we have

$$\lim_{\eta \rightarrow 0} \frac{1}{z + i\eta} = \text{P.V.} \left(\frac{1}{z} \right) - i\pi\delta(z),$$

$$\sum_{\alpha} \delta(\varepsilon - \varepsilon_{\alpha,s}) |\psi_{\alpha,s}(\mathbf{r})|^2 = -\frac{1}{\pi} \text{Im} \lim_{\eta \rightarrow 0} \sum_{\alpha} \frac{|\psi_{\alpha,s}(\mathbf{r})|^2}{\varepsilon - \varepsilon_{\alpha,s} + i\eta} = -\frac{1}{\pi} \text{Im} G_s(\mathbf{r}, \mathbf{r}; \varepsilon)$$

Therefore, in terms of the Green function, the electron density is given by

$$G_s(\mathbf{r}, \mathbf{r}'; \varepsilon) = \lim_{\eta \rightarrow 0} \sum_{\alpha} \frac{\psi_{\alpha,s}^*(\mathbf{r}) \psi_{\alpha,s}(\mathbf{r}')}{\varepsilon - \varepsilon_{\alpha,s} + i\eta}$$

$$\rho_s(\mathbf{r}) = -\frac{1}{\pi} \text{Im} \int_{-\infty}^{\infty} f_T(\varepsilon - \varepsilon_F) G(\mathbf{r}, \mathbf{r}; \varepsilon) d\varepsilon. \quad \text{At } T = 0, \text{ we have: } \rho_s(\mathbf{r}) = -\frac{1}{\pi} \text{Im} \int_{-\infty}^{\varepsilon_F} G_s(\mathbf{r}, \mathbf{r}; \varepsilon) d\varepsilon$$

Electronic density of state and Green function

The density of states is given by

$$\rho_s(\varepsilon) = \sum_{\alpha} \delta(\varepsilon - \varepsilon_{\alpha,s}) \int_{\infty} |\psi_{\alpha,s}(\mathbf{r})|^2 d^3\mathbf{r} = \int_{\infty} \left[\sum_{\alpha} \delta(\varepsilon - \varepsilon_{\alpha,s}) |\psi_{\alpha,s}(\mathbf{r})|^2 \right] d^3\mathbf{r}$$

Applying the relation shown from the previous slide

$$\sum_{\alpha} \delta(\varepsilon - \varepsilon_{\alpha,s}) |\psi_{\alpha,s}(\mathbf{r})|^2 = -\frac{1}{\pi} \text{Im} G_s(\mathbf{r}, \mathbf{r}; \varepsilon)$$

we have

$$\rho_s(\varepsilon) = -\frac{1}{\pi} \text{Im} \int_{\infty} G_s(\mathbf{r}, \mathbf{r}; \varepsilon) d^3\mathbf{r} = -\frac{1}{\pi} \text{Im} \sum_n \int_{\Omega_n} G_s(\mathbf{r}, \mathbf{r}; \varepsilon) d^3\mathbf{r}$$

The density of states in space domain Ω_n is given by

$$\rho_s^n(\varepsilon) = -\frac{1}{\pi} \text{Im} \int_{\Omega_n} G_s(\mathbf{r}, \mathbf{r}; \varepsilon) d^3\mathbf{r}$$

Contour integration

The electron density calculation can be carried out along a contour in the complex plane

$$\begin{aligned}\rho_s(\mathbf{r}) &= -\frac{1}{\pi} \text{Im} \int_{-\infty}^{\varepsilon_F} G_s(\mathbf{r}, \mathbf{r}; \varepsilon) d\varepsilon \\ &= \sum_c \rho_{c,s}(\mathbf{r}) - \frac{1}{\pi} \text{Im} \int_{\varepsilon_b}^{\varepsilon_F} G_s(\mathbf{r}, \mathbf{r}; \varepsilon) d\varepsilon \\ &= \sum_c \rho_{c,s}(\mathbf{r}) - \frac{1}{\pi} \text{Im} \oint_{\varepsilon_b}^{\varepsilon_F} G_s(\mathbf{r}, \mathbf{r}; z) dz\end{aligned}$$

$\rho_{c,s}(\mathbf{r})$ is the electron density of core state c with spin s .

$$\rho_v(\varepsilon) = -\frac{1}{\pi} \text{Im} \sum_{s=\uparrow,\downarrow} \sum_n \int_{\Omega_n} G_s(\mathbf{r}, \mathbf{r}; \varepsilon) d^3\mathbf{r}$$

with $\varepsilon_b \leq \varepsilon \leq \varepsilon_F$.

So is the band energy calculation,

$$\begin{aligned}\int_{\varepsilon_b}^{\varepsilon_F} \varepsilon \rho_v(\varepsilon) d\varepsilon &= -\frac{1}{\pi} \text{Im} \sum_{s=\uparrow,\downarrow} \sum_n \int_{\varepsilon_b}^{\varepsilon_F} \varepsilon \left(\int_{\Omega_n} G_s(\mathbf{r}, \mathbf{r}; \varepsilon) d^3\mathbf{r} \right) d\varepsilon \\ &= -\frac{1}{\pi} \text{Im} \sum_{s=\uparrow,\downarrow} \sum_n \oint_{\varepsilon_b}^{\varepsilon_F} z \left(\int_{\Omega_n} G_s(\mathbf{r}, \mathbf{r}; z) d^3\mathbf{r} \right) dz\end{aligned}$$